Electrical Drive Inductive Coupling

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Abstract The paper presents a computer analysis of inductive coupling of the electromagnetic compatibility (EMC) problem. Its focus is on power electronics and electrical drives and tests performed by a numerical computer simulation that can disclose suite surprising findings about EMC problems.

Keywords: electromagnetic compatibility, power electronics, converters, inverters.

1 Introduction

Importance of electromagnetic compatibility (EMC) of all electrical products has been rapidly growing during the last decade. The environment is increasingly polluted by electromagnetic energy. The interference impact on the surroundings is being doubled every three years and covers a large frequency range.

 Equipment disturbances and errors have become more serious as a consequence of the growth of the electronic circuit complexity. According to new technical legislation and also economic consequences, the EMC concept of all products must be strictly observed [1]. It must start with the specification of the equipment performance and end with the equipment installation procedures.

2 EMC and environmental waste

We all know the environmental pollution problems caused by solid, liquid and gaseous wastes. We are aware of most of these pollutants through our senses.

Due to the increasing life standard, contamination of our environment by the electromagnetic energy is constantly increasing too. Since human beings have no organs for perception of such contamination, they cannot perceive it. The great producers of such waste are electronic systems developed by man and meant to be effective within these electromagnetic surroundings producing, of course, electromagnetic waste in turn [2].

 On one side, interferences are deliberately or involuntarily produced. The place of their origin is called interference source. On the other side, devices may be hindered in their function by such interferences. Those objects are called interference objects.

 The possible interfaces between sources and objects are shown in Fig. 1. There are four basic types of coupling that can realize these interfaces.

Fig. 1. Interference diagram

3 EMC – the interference mechanism

The interference mechanism can be described in a simplified form as follows. The interference source can be for instance a power semiconductor converter or motor. Interference is produced in the interference source getting into electronics in undesirable ways and is due to various effects distorting signals. Transmission can be direct, for example by galvanic coupling between interference source and interference sink. Interference can be spread through air or via ducts, or coupled *Received 8. March, 2005* inductively or capacitively into signal lines [3].

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Development of power semiconductor elements has caused vehement evolution of the power electronics branch in the last ten years. To investigate the converter functionality, it was necessary first to theoretically analyze and then practically verify its assumed activity. Now, we can eliminate the laborious theoretical analysis by a numerical computer simulation, which can disclose quite surprising findings about EMC [4].

4 Inductive coupling

Inductive coupling is typical for two and more galvanically separated electric loops at the moment when the smaller one is driven by a time variablecurrent creating the corresponding, time-variable magnetic field [5]. In such case their mutual intercircuit effect is expressed as a function of the slope of the current increase or decrease, circuit environmental magnetic property as well as circuit geometric dimensions.

 To predict the intercircuit inductive coupling, our focus will be on two electric loops l_1 and l_2 with currents i_l and i_2 . We will try to determine the effect of loop l_l on loop l_2 (Fig. 2).

Fig. 2. Investigated loops

According to the Maxwell's equation for a quasistationary magnetic field

$$
rot\ \overline{E} = -\frac{\partial \overline{B}}{\partial t} \tag{1}
$$

and following its integral form

$$
\int_{S} rot \ \overline{E} .d\overline{S} = -\int_{S} \frac{\partial \overline{B}}{\partial t} .d\overline{S} = -\frac{\partial}{\partial t} \int_{S} \overline{B} .d\overline{S}
$$
 (2)

and after applying the Stoke's theorem, we obtain the equation for the induced voltage [6];

$$
u_{i2} = -N \cdot \frac{\partial \phi_1}{\partial t} = -\frac{\partial \psi_1}{\partial t} = -M \frac{\partial i_1}{\partial t}
$$
 (3)

where *M* is the coefficient of the mutual inductance. For the magnetic flux Ψ definition the equation

$$
\phi_1 = \oint_{l_2} \overline{A}_2 \cdot d\overline{l}_2 \tag{4}
$$

is valid where \overline{A}_2 is the vector of the magnetic field potential created by the current i_l . We can calculate the value of this vector by the following equation:

$$
\overline{A}_2 = \frac{\mu i_1}{4\pi} \oint_{l_1} \frac{d\overline{l}_1}{r_{12}}.
$$
\n(5)

After substituting the last equation with the equation valid for the magnetic flux ϕ_l , the next relation is obtained:

$$
\phi_1 = \oint_{l_2} \left[\frac{\mu i_1}{4\pi} \oint_{l_1} \frac{d\bar{l}_1}{r_{12}} \right] d\bar{l}_2 = \frac{\mu i_1}{4\pi} \oint_{l_1} \oint_{l_2} \frac{d\bar{l}_1 d\bar{l}_2}{r_{12}} \tag{6}
$$

and then

$$
u_{i2} = -\frac{\partial \left(\frac{\mu i_1}{4\pi} \oint_{i_1} \oint_{i_2} \frac{d\bar{l}_1 d\bar{l}_2}{r_{i2}}\right)}{\partial t} = -\frac{\left(\frac{\mu}{4\pi} \oint_{i_1} \oint_{i_2} \frac{d\bar{l}_1 d\bar{l}_2}{r_{i2}}\right) \partial i_1}{\partial t} = -M \frac{\partial i_1}{\partial t}.
$$
\n(7)

For the practical use, it is more advantageous to express the induced voltage in the form of a differential:

$$
u_{i} = -\frac{di}{dt} \cdot \sum_{i=1}^{m} \sum_{j=1}^{k} \frac{\mu}{4\pi} \frac{dl_{1i}.dl_{2j}. \cos \gamma_{dlij}}{r_{ij}}.
$$
 (8)

If we know the geometrical dimensions of the investigated loops (Fig. 3) and want to determine their mutual inductive coupling then we can use the next relation (9) for the induced voltage. It is based on the 3D Cartesian coordinate system.

Fig. 3. Geometric dimensions of the investigated loops

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$$
u_{i} = \frac{di}{dt} \sum_{i=1, j=1}^{m} \frac{\mu}{4\pi} \frac{(A_{x2} - A_{x1}) \cdot (B_{x2j} - B_{x1j}) + (A_{y2} - A_{y1}) \cdot (B_{y2j} - B_{y1j}) + (A_{z2i} - A_{z1}) \cdot (B_{z2j} - B_{z1j})}{\sqrt{\left(\left(B_{x1j} + \frac{|B_{x2j} - B_{x1j}|}{2} \right) - \left(A_{x1j} + \frac{|A_{x2j} - A_{x1}|}{2} \right) \right)^{2} + \left(\left(B_{y1j} + \frac{|B_{y2j} - B_{y1}|}{2} \right) - \left(A_{y1} + \frac{|A_{y2} - A_{y1}|}{2} \right) \right)^{2} + \left(\left(B_{z1j} + \frac{|B_{z2j} - B_{z1j}|}{2} \right) - \left(A_{z1j} + \frac{|A_{y2} - A_{y1}|}{2} \right) \right)^{2}} \right)}
$$
(9)

For a global solution of the inductive coupling part of the EMC problem inside the overall electric power system, it is necessary to analyze the circuit globally paying due regard to the mutual intercircuit inductance coupling. The result is the following integral-differential system of equations:

$$
u_{cc1} = R_{c1} \dot{u}_1 + L_{c1} \frac{di_1}{dt} + \frac{1}{C_{c1}} \int \dot{u}_1 dt + \sum_{\substack{j=1 \ j \neq 1}}^k u_{ij}
$$
 (10)

$$
\vdots
$$
\n
$$
u_{cck} = R_{ck} \cdot i_k + L_{ck} \cdot \frac{di_k}{dt} + \frac{1}{C_{ck}} \int i_k \cdot dt + \sum_{\substack{j=1 \ j \neq k}}^k u_{ij} \tag{11}
$$

For this purpose it is very suitable to explore the existing simulation programs such as for instance the PSPICE program utilized worldwide.

 In the next part, we will try to determine the effect of the one-quadrant impulse converter on the sensing circuit as it shown in Fig. 4. The circuit dimensions are $a = 0.2$ m, $b = 0.3$ m, $c = 0.1$ m, $d = 0.05$ m, $e = 0.005$ m. The radius of the copper wires is $R = 0.0006$ m and the relative permitivity of the circuit environment is μ_r = 0.991.

Fig. 4. Investigated circuit

The inductance of the first loop is given as

$$
L_1 = L_{e1} + L_{i1} = \frac{\mu_0 b}{\pi} \ln \frac{a - R}{R} + \frac{\mu_0 a}{\pi} \ln \frac{b - R}{R} + \frac{\mu 2(a + b)}{8\pi} = 1.294 \mu H
$$
\n(12)

and of the second as

$$
L_2 = L_{e2} + L_{r2} = \frac{\mu_0 c}{\pi} \ln \frac{d - R}{R} + \frac{\mu_0 d}{\pi} \ln \frac{c - R}{R} + \frac{\mu 2(c + d)}{8\pi} = 0.294 \mu H
$$
\n(13)

The mutual inductance *M* calculated from the above mentioned equation is $M = 477.4$ nH. The magnetic coupling coefficient *k* is given as

$$
k = \frac{M}{\sqrt{L_1 + L_2}} = 0.774\,. \tag{14}
$$

Now we can use the PSPICE simulation program for solving the inductive coupling problem between the two circuits [7]. Parameters of the circuit simulation are R_Z = 11.66 Ω , L_Z = 400 µH, R = 10 Ω , R_G = 100 Ω and U_{CC} = 70V. The schematic connection is shown in Fig. 5. The IGBT transistor Q was switched on at the frequency 10 kHz and the switch on/off ratio was 0.5.

Fig. 5. Simulation circuit

Simulation results are shown in Fig. 6. Results obtained with measurements are shown in Figs. 7 and 8 and switching details in Figs. 9 and 10.

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50V 10mV 20µ**s SAVE**

0

 i_c $2A/d$

Fig. 9. Switching on voltage u_{CE} and current i_C

Fig. 10. Switching off voltage u_{CE} and current i_C

A comparison of the simulated and measured results shows that peaks of transistor current ic have the same values, i.e. 8.4 A, in both cases. The same values, i.e. 4.4 A, have both the simulated and measured transistor current at the moment when transistor is switched off. There is a small difference only between the simulated and measured curves of the transistor voltage uCE. The overvoltage generated at the transistor switching off reaches the value of 150 V for the simulated result. However, the corresponding overvoltage has only the value of 130 V for the measured result. Peaks of the simulated and measured induced voltages have the same values of Ui1 = -2.2V, Ui2 = $5.02V$, Ui1 = 2.1V. This means that such method is acceptable for inductive coupling investigation of the EMC problem.

 To improve the obtained results, the numerical solution of the magnetic field by finite element method program was also used. The result of such analysis is shown in Fig. 11.

 From the "integral result" data window it is seen, that the value of the magnetic flux inside the sensing circuit is $3.317.10^{-9}$ Wb. Based on the basic program property allowing semi-real 3D space simulation with the $3rd$ dimension equal only to the basic unit of the depth (1mm), we multiplied the obtained value of the magnetic flux by the value of the sensing circuit depth c $= 100$ mm. The total magnetic flux was then $331.7.10^{-9}$ Wb. This flux was excited by the peak circuit current 8.4 A, the rising time of which was 120 ns. On the basis of the above equations, the first peak of the induced voltage can be calculated as

$$
U_{ip1} = \frac{\Delta \Phi_1}{\Delta t_1} = \frac{0 - 331.7.10^{-9}}{140.10^{-9}} = \frac{-331,7.10^{-9}}{140.10^{-9}} = -2.369 \text{ V}.
$$
\n(15)

Similarly, it is possible to calculate the rest of the peaks of the induced voltage *ui*:

$$
U_{ip2} = \frac{\Delta \Phi_2}{\Delta t_2} = \frac{331.7 \cdot 10^{-9} - 55.3 \cdot 10^{-9}}{55.10^{-9}} = \frac{2764.10^{-9}}{55.10^{-9}} = 5.025 \text{V}
$$
\n
$$
U_{ip3} = \frac{\Delta \Phi_3}{\Delta t_3} = \frac{173.7 \cdot 10^{-9} - 0}{80.10^{-9}} = \frac{173.7 \cdot 10^{-9}}{80.10^{-9}} = 2.171 \text{ V}.
$$
\n(17)

The results obtained by the finite element numerical simulation method are again confirming the correctness of the above mentioned methods.

5 Conclusion

The performed analyses indicate that the fast power field effect transistor switching can produce the induced voltage with the value of some volts up to some tenths' of volts in the nearby circuits. It is also evident that the magnitude of the induced voltage depends on the magnetic flux slope. This means that fast switching of small currents can generate large peaks of the induced voltage, too.

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Fig. 11. The finite element simulation method of the magnetic field

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